# SHORTER COMMUNICATIONS

# A NOTE ON THE LOCKHART-MARTINELLI CORRELATION IN THE TURBULENT-TURBULENT CASE

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#### NOMENCLATURE

|    | <b>.</b> |           |
|----|----------|-----------|
| b. | Blasius  | exponent: |

K, dimensionless parameter, equation (5);

m, dimensionless parameter to be fitted to the

experimental data;

X<sub>TT</sub>, Lockhart-Martinelli parameter for the turbulent-turbulent case.

## Greek symbols

α, void fraction;

y, slip ratio;

 $\mu_{\rm g}$ , vapor viscosity;

 $\mu_{\rm b}$ , liquid viscosity;  $\rho_{\rm w}$ , vapor density;

 $\rho_{\mathfrak{g}},$  vapor density;  $\rho_{\mathfrak{h}},$  liquid density;

 $\phi_{\rm L}$ , parameter related to the void fraction  $\alpha$  by

equation (2); χ, vapor quality.

For two phase turbulent-turbulent flow in pipes, the void fraction  $\alpha$  depends on the vapor quality,  $\chi$ , on the vapor and liquid densities  $\rho_{\mathbf{r}}$ ,  $\rho_{\mathbf{l}}$  on the vapor and liquid viscosities  $\mu_{\mathbf{r}}$ ,  $\mu_{\mathbf{l}}$ . That  $\alpha$  actually only depends on the parameter  $X_{\mathbf{T}\mathbf{l}}$  is the physical content of the Lockhart-Martinelli correlation [1],  $X_{\mathbf{T}\mathbf{l}}$  being defined by

$$X_{TT} = [(1 - \chi)/\chi]^{(2-b)/2} (\rho_g/\rho_l)^{1/2} (\mu_l/\mu_g)^{(b/2)}$$
 (1)

where b denotes the Blasius exponent.

The curve  $\alpha(X_{TT})$  is usually fitted by the equations

$$\alpha = 1 - (1/\phi_L),$$

$$\phi_L^2 = 1 + 21/X_{TT} + 1/X_{TT}^2.$$
(2)

Using such a fit is rather troublesome when performing analytical calculations. Moreover equation (2) yields an unphysical singularity for  $\partial \alpha/\partial \chi$  as  $\chi \to 0$ . As a substitute for equation (2) we propose the relation

$$\alpha = 1 - (m/2) X_{11}^{2(2-b)} - \{ [1 + (m/2) X_{11}^{2(2-b)}]^2 - 1 \}^{1/2}$$
(3)

where m denotes a dimensionless parameter to be fitted to the experimental data.

A comparison of the fit of equations (2) and (3) is given in Table 1, with m=0.054. The agreement is quite good, except for  $\chi \to 0$ , since equation (3) does not generate any unphysical singularity.

 $\alpha(\chi)$ , as defined by equation (3), only depends on  $X_{\rm TT}$ , according to the Lockhart-Martinelli correlation. This is the only case of equation (3) which can be turned into the simple analytical relation

$$\chi = (\alpha mK)/[(1-\alpha)^2 + \alpha mK] \tag{4}$$

where K is defined by

$$K = [(\rho_{\rm g}/\rho_{\rm l})^{1/2} (\mu_{\rm l}/\mu_{\rm g})^{b/2}]^{2/(2-b)}.$$
 (5)

Let us now consider the slip ratio y, which is known to satisfy the identity

$$\gamma = (\rho_1/\rho_2) \left[ (1-\alpha)/\alpha \right] \left[ \chi/(1-\chi) \right]. \tag{6}$$

Using equation (4), this simplifies to

$$\gamma = (\rho_{\text{V}}/\rho_{\text{x}})[mK/(1-\alpha)]. \tag{7}$$

Some properties of the correlations, equations (4) and (7) are:

- (1) From the numerical point of view, the agreement with the usual correlation, equation (2), is quite good.
  - (2) They do not yield unphysical singularities as  $\chi \to 0$ .
- (3) Their analytical form is very simple.
- (4) They only depend on the parameter m, which must be

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Table 1. Comparison of the correlations equations (2) and (3)

| $X_{TT}$ | χ                  | α[equation (2)] | α[equation (3)] |
|----------|--------------------|-----------------|-----------------|
| 55.21    | 2×10 <sup>-4</sup> | 0.149           | 0.137           |
| 30.20    | $4 \times 10^{-4}$ | 0.232           | 0.223           |
| 21.22    | $6 \times 10^{-4}$ | 0.291           | 0.284           |
| 13.52    | $8 \times 10^{-4}$ | 0.337           | 0.330           |
| 16.60    | $1 \times 10^{-3}$ | 0.374           | 0.368           |
| 7.436    | $2 \times 10^{-3}$ | 0.490           | 0.487           |
| 5.221    | $3 \times 10^{-3}$ | 0.555           | 0.553           |
| 4.062    | $4 \times 10^{-3}$ | 0.599           | 0.598           |
| 3.342    | $5 \times 10^{-3}$ | 0.632           | 0.631           |
| 2.849    | $6 \times 10^{-3}$ | 0.657           | 0.656           |
| 2.489    | $7 \times 10^{-3}$ | 0.677           | 0.677           |
| 2.214    | $8 \times 10^{-3}$ | 0.694           | 0.694           |
| 1.997    | $9 \times 10^{-3}$ | 0.708           | 0.709           |
| 1.821    | $1 \times 10^{-2}$ | 0.721           | 0.721           |
| 0.9873   | $2 \times 10^{-2}$ | 0.793           | 0.794           |
| 0.6877   | $3 \times 10^{-2}$ | 0.828           | 0.829           |
| 0.5306   | $4 \times 10^{-2}$ | 0.849           | 0.851           |
| 0.4330   | $5 \times 10^{-2}$ | 0.865           | 0.866           |
| 0.3661   | $6 \times 10^{-2}$ | 0.877           | 0.878           |
| 0.3172   | $7 \times 10^{-2}$ | 0.886           | 0.887           |
| 0.2798   | $8 \times 10^{-2}$ | 0.894           | 0.894           |
| 0.2501   | $9 \times 10^{-2}$ | 0.900           | 0.901           |
| 0.2260   | $1 \times 10^{-1}$ | 0.906           | 0.906           |
| 0.08691  | 0.25               | 0.948           | 0.945           |
| 0.03342  | 0.50               | 0.974           | 0.968           |
| 0        | 1                  | 1               | 1               |

In order to relate the vapor quality  $\chi$  and the parameter  $X_{\rm TT}$  the following parameters have been used (corresponding to sodium):  $\rho_{\rm I}=0.7217\,{\rm g\,cm^{-3}},\,\rho_{\rm g}=0.48.10^{-3}\,{\rm g\,cm^{-3}},\,\mu_{\rm I}=0.1475$  centipoises,  $\mu_{\rm g}=0.02009$  centipoises, b=0.26. In equation (3) the parameter m has been fitted to m=0.054.

fitted to the experimental results. The usual correlation, equation (2), could be adjusted either by modifying the coefficients of the series in  $1/X_{\rm TT}$  or by adding to it further

### REFERENCES

1. R. W. Lockhart and R. C. Martinelli, Chem. Engng Prog. 45, 39 (1949).

 $\label{eq:central_loss} Acknowledgements — I \ wish \ to \ thank \ M. \ Ph. \ Berna, \ CEA/DSN/Cadarache for illuminating discussions.$